

1 **Hurricanes on Tidally Locked Terrestrial Planets**

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7
8 **Abstract: Are there hurricanes on exoplanets? Tidally locked terrestrial planets around**
9 **M dwarfs are the main targets of space missions for finding habitable exoplanets. Whether**
10 **hurricanes can form on this kind of planet is important for their climate and habitability.**
11 **Using a high-resolution global atmospheric circulation model, we show that hurricanes can**
12 **form on the planets but not on all of them. For planets near the inner edge of the habitable**
13 **zone of late M dwarfs, there are more and stronger hurricanes on both day and night sides.**
14 **Hurricane theories on Earth are applicable when atmospheric compositions are similar to**
15 **Earth. However, if background atmosphere is lighter than H₂O, hurricanes can hardly be**
16 **produced because convection is always inhibited due to the effect of mean molecular weight,**
17 **similar to that on Saturn. These results have broad implications on the precipitation, ocean**
18 **mixing, climate, and atmospheric characterization of tidally locked planets.**

23 **Introduction**

24 Hurricanes (also known as tropical cyclones or typhoons) are low-pressure weather systems with
25 well-organized convection and are one of the most destructive disasters on Earth because they can
26 induce strong winds, rapid increases in local sea level, and heavy precipitations (1,2). Hurricanes
27 can also enhance the vertical mixing of heat and nutrients in the ocean, increase horizontal oceanic
28 heat transport, and subsequently influence global climate (3-5). For instance, the curl of the strong
29 winds can cause divergence and convergence in the upper ocean, producing regions of up-welling
30 and down-welling and enhancing the exchanges between surface and subsurface oceans. Therefore,
31 it is important and interesting to know whether hurricanes can form on other potentially habitable
32 planets beyond Earth.

33 In this study, we focus on tidally locked terrestrial planets around M dwarfs, due to their
34 relatively large planet-to-star ratios and frequent transits in observation. These planets differ from
35 Earth in three major aspects: the uneven distribution of stellar energy between the permanent day
36 and night sides, the slow rotation rate due to strong tidal force, and the redder stellar spectrum than
37 the Sun. Several atmospheric general circulation models (AGCMs) have been employed and
38 modified to simulate and understand atmospheric and climatic dynamics of the planets (6-24).
39 These AGCMs have horizontal resolutions that are always equal to or larger than 300 km, so that
40 most of their studies focus on planetary-scale phenomena, such as global-scale Walker circulation,
41 equatorial superrotation, and forced Rossby and Kelvin waves. No work has investigated synoptic
42 phenomena except Bin et al. (25). Based on the output data of an AGCM, they estimated genesis
43 potential index of hurricanes and showed that the probability of hurricane formation is low for
44 planets in the middle range of the habitable zone of M dwarfs. However, model resolution they
45 used was not able to directly simulate hurricanes and whether the empirical index can be applied to
46 exoplanets was unknown.

47 In this study, we explicitly simulate hurricane formation on tidally locked terrestrial planets
48 with a high-resolution AGCM (see Materials and Methods). Horizontal resolution of the model is
49 ≈ 50 km, under which the cyclogenesis, intensity, and structure of large hurricanes can be well
50 resolved. The atmosphere is set to be dominated by N₂, O₂, CO₂, H₂, or He with a mean surface air
51 pressure of 1 bar. Four different planetary rotation periods have been examined, 6, 10, 20, and 40
52 earth days. The surface is covered by seawater everywhere without land. Due to computation time
53 limits, surface air temperatures are specified, but their magnitude and spatial pattern are set to be
54 close to the results of relatively low-resolution AGCM simulations those have a slab ocean. We
55 find that hurricanes can indeed form on the tidally locked planets but not on all of them.

56

57 **Results**

58 **Hurricanes on tidally locked planets**

59 Figure 1 shows the results of the control experiment for an aquaplanet orbiting close to the inner
60 edge of the habitable zone around a late M dwarf of 2600 K. The rotation period is set to 6 earth
61 days. Both the day and night sides are set to warm, and the day-to-night surface temperature contrast
62 is small (see Fig. 2A). This experimental design is due to the fact that the day-to-night atmospheric
63 latent heat transport is very efficient for planets near the inner edge (9,10). The high surface
64 temperature, weak day-to-night contrast, and relatively fast rotation rate (comparing to planets in
65 the middle range of the habitable zone or planets orbiting around hotter stars) benefit hurricane
66 formation.

67 Clearly, there are hurricanes in the control experiment. In the mature stage of the hurricanes,
68 maximum wind speed reaches ≈ 30 -50 m s⁻¹ (Fig. 1A), surface air pressure at the center is ≈ 950 -
69 980 hPa (Fig. 1B), precipitation reaches as high as 200-500 mm per day due to strong convection
70 especially near the eyewall (Fig. 1C), and relative vorticity near the surface is in the order of 10⁻⁴
71 s⁻¹ (Fig. 1D), close to those on Earth (1,2). The surface wind speed increases as the eyewall is

72 approached from outside, but inside the eyewall the winds as well as precipitation weaken rapidly.
73 The winds rotate counter-clockwise in the northern hemisphere and clockwise in the southern
74 hemisphere due to the Coriolis force although in this experiment it is much smaller than that on
75 Earth. The precipitation exhibits well-defined spiral bands rather than uniformly distributed
76 throughout the region of the hurricane. The clear patterns of eye–eyewall and spiral rain bands
77 suggest that the model resolution of 50 km is good in resolving the hurricanes.

78 For vertical cross structure (Fig. S1), the tangential wind component dominates the flow
79 through the system although the radial wind is also significant. In the boundary layer, the winds
80 flow toward the low-pressure center. Within the hurricane, upward motion is robust and tilts radially
81 outward. The maximum ascendance near the surface locates at a distance of ≈ 250 km outward from
82 the hurricane center rather than at the center itself. Indeed, weak downward motion comprises the
83 center. Due to the ascendance and deep convection, relative humidity is high in the hurricane. Latent
84 heat release from the convection and adiabatic warming by compression from the subsidence in the
85 eye produce a warmer region of air with temperatures of ≈ 4 K above the environmental value. This
86 can be called as a “warm core”, which is one of the most characteristic features of hurricane. The
87 warm core can also be viewed from the equivalent potential temperature anomaly. Hurricanes on
88 the night side (Fig. S2) are smaller in horizontal size compared to that on the day side, ≈ 500 versus
89 1500 km, but the vertical structures are similar.

90 Statistical analysis shows that in the control experiment there are four preferred regions for
91 hurricane genesis: the northern and southern tropics of the day side near the substellar point and the
92 middle-to-high latitudes of the night side on each hemisphere (black dots in Fig. 2A). Hurricane
93 formation is largely determined by small-scale convection, large-scale environmental condition,
94 and the interactions between them. Below, we explore the underlying mechanisms through two
95 ways.

96 One way is the positive feedback between cumulus convection and larger-scale disturbance,
97 known as the Conditional Instability of the Second Kind (CISK (26-29)). On tidally locked planets,
98 long-term mean atmospheric circulation is characterized by large-scale Rossby waves on the west
99 and pole of the substellar point and Kelvin waves on the east of the substellar point (Fig. 2B),
100 excited from the uneven stellar radiation distribution (6,31). The wave pattern is similar to the
101 tropical Matsuno-Gill pattern on Earth (30,32), but the meridional (south-north) scale is larger,
102 $\approx 10,000$ versus 3,000 km. The Rossby waves have one low-pressure center on each hemisphere,
103 whose corresponding environmental vertical motion is updrafts and relative vorticity is positive
104 (negative) on the northern (southern) hemisphere. This low pressure system favors the onset of the
105 CISK feedback: Surface winds spiral into the low-pressure center and create horizontal
106 convergence; this low-level convergence enhances relative vorticity through vortex stretching,
107 increases upward motion following the conservation of mass, and meanwhile brings water vapor
108 into the center, amplifying cumulus convection and release of latent heat; the latent heat release
109 warms the air and lowers the air density through forcing more upper-level air to move outward
110 away from the center, subsequently reducing the surface pressure; the lower surface pressure further
111 enhances the low-level convergence and increases the growth rate of the relative vorticity through
112 vortex stretching (Fig. 2D-F). This feedback is the key in promoting the growth of small-scale
113 disturbances to hurricanes in the background low-pressure regions. It is similar to that on Earth:
114 hurricane generally forms in the monsoon troughs and the confluence zones where the surface
115 pressure is relatively low, collocated with high cyclonic vorticity, convergent surface winds, and
116 divergent winds aloft (1,2). Moreover, during the formation phase, the latent heat flux from the
117 surface to the boundary layer increases strongly (Fig. 2G), which also contributes to intensify the
118 hurricane through the feedback of wind-induced surface energy exchange (WISHE (2,28)).

119 The lifetime of the hurricanes on the day side is ≈ 40 -50 earth days, longer than that on Earth.
120 This is mainly due to the absence of continents and the warm surface everywhere in the experiment.
121 On the night side, the lifetime is shorter, ≈ 10 -20 earth days.

122 Another way is the empirical equation of genesis potential index (GPI (25,33,34)). The index
123 combines five environmental factors to predict the potential of hurricane formation, including
124 planetary vorticity, relative vorticity, relative humidity, potential intensity, and wind shear (see
125 Methods). Comparison between Figs. 2A and 3A reveals a positive correlation between the location
126 of hurricane genesis and large values of GPI. In the four hurricane regions, GPI values are large
127 because the relative vorticity is great, relative humidity is high, potential intensity is large, and/or
128 vertical wind shear is weak (Fig. 3B-F). These favor hurricane formations there. For example, when
129 the shear is strong, an initial disturbance will be ventilated by cooler or drier air and thereby
130 temperature and moisture anomalies are hard to maintain (35). In this experiment, the vertical wind
131 shear is strong especially in the tropics of the night side associated with atmospheric superrotation
132 (6,7) and in the extratropics of the day side (Fig. 3F), so that there is nearly no hurricane formation
133 there. Note that the applicability of the empirical GPI index on Earth to the tidally locked planet is
134 mainly due to that we employed an earth-like atmosphere here; when the atmospheric composition
135 is quite different from Earth, GPI is not a good index as addressed below.

136 The formation of hurricanes on the night side is surprising because the night side has no stellar
137 radiation and the long-term mean vertical motion is downwelling rather than upwelling. In this
138 experiment of planet close to the inner edge of the habitable zone, however, the night-side surface
139 is warm, the surface temperature gradient is small (Fig. 2A), there are a few short-time small low-
140 pressure regions (Fig. 2C), the planetary vorticity is relatively high (Fig. 3B), and the vertical wind
141 shear is weak (Fig. 3F), at the middle-to-high latitudes. These factors promote hurricane genesis
142 there. But, when the surface temperature is decreased or the rotation rate is slowed down, there is
143 fewer or no hurricane on the night side (see below).

Effects of planetary rotation rate and surface temperature

In order to test the effect of rotation rate, we performed three experiments in which rotation period is increased (i.e., the rotation rate is decreased) while other experimental designs are left to be the same as that in the control experiment. In the case of 10 days, the hurricane frequency does not change much on the day side but decreases significantly on the night side (Fig. 4A). In the case of 20 days, there is nearly no hurricane on the night side and the number of hurricanes on the day side also decreases largely (Fig. 4B). For the case of 40 days, hurricane only forms at regions very close to the substellar point (Fig. 4C). This trend as a function of rotation period is mainly due to three factors: the direct weakness of planetary vorticity, the reduction of relative vorticity due to that the atmosphere becomes so steady that waves and disturbances become less active, and the decrease of relative humidity on the night side (but increase on the day side) due to the strengthening of the thermal-driven global Walker circulation (Fig. S5).

When the surface temperature is decreased, hurricane is also harder to form. As the maximum surface temperature is set to 308 K and the night-side surface temperature is set to 275 K, fewer hurricanes form in the vicinity of the substellar point and no hurricane on the night side (Fig. 4D). When the maximum surface temperature is set to 301 K and the night-side surface temperature is set to 268 K, there are only two hurricane events during the integration of four earth years (Fig. 4E), consistent with the GPI prediction in Bin et al. (25). Note that the temperature of 301 K is close to the tropical surface temperatures on Earth. This suggests hurricane formation on tidally locked planets requires a warmer surface, due to their slower rotation rates and stronger wind shears. For planets with both slow rotation and low temperature, no hurricane can form (Fig. 4F). The decreasing trend of hurricane formation as a function of reduced surface temperature is due to two main processes: the relative humidity and potential intensity decrease because of the cooler surface and weaker upwelling and convection, and the vertical wind shear becomes much stronger due to the enhanced temperature gradients between day and night sides (Fig. S6). Moreover, when the

169 night-side surface temperature is low, air convergence from the night side to the day side brings
170 cool air rather than warm air into the substellar region, suppressing hurricane formation there.

172 **Effect of bulk atmospheric composition**

173 Atmospheric compositions on terrestrial exoplanets are unknown yet. Here we do a preliminary
174 investigation of how atmospheric molecular weight influences hurricane formation under a uniform
175 surface temperature of 301 K (see Methods). When the background atmosphere is set to H₂ or He,
176 there is no hurricane, in contrast to the experiments of N₂, O₂, and CO₂ (Fig. 5), although GPI value
177 is comparable to or even larger than that shown above. This is due to the fact that the condensate–
178 H₂O is heavier than H₂ and He, so that any disturbance that brings water vapor upward will cause
179 the density of a moist parcel to be larger than its surrounding environment, similar to the condition
180 in Saturn’s atmosphere (36-38). This process induces a negative buoyancy and stabilizes the
181 atmosphere against convection. Simply, it can be understood using the ideal gas equation $p = \rho R_d T_v$
182 and the virtual temperature of

$$183 \quad T_v = \frac{p}{p + (\epsilon - 1)e} T, \quad (1)$$

184 where p is total air pressure, e is partial pressure of the condensate, ρ is air density, R_d is the gas
185 constant of dry air, ϵ is the molecular weight ratio of water vapor to the dry air, and T is air
186 temperature. For a H₂-dominated (or He-dominated) atmosphere, ϵ is equal to 9 (or 4.5), so that T_v
187 is smaller than T . Therefore, a moist parcel is heavier than a dry parcel under the same p and T , and
188 moist convection is inhibited, which is opposite to the condition on Earth. Moreover, in the
189 experiments of N₂, O₂, and CO₂, a clear trend is that the hurricane size decreases as the mean
190 molecular weight is increased. This is due to the fact that atmospheric scale height is inversely
191 proportional to the mean molecular weight and subsequently the Rossby deformation radius (see
192 Materials and Methods) becomes smaller.

194 **Discussion**

195 We find that hurricanes can form on tidally locked planets especially for those orbiting near the
196 inner edge of the habitable zone of late M dwarfs. For planets in the middle range of the habitable
197 zone, hurricanes are relatively fewer. Storm theories of Earth and Saturn can be used to understand
198 the hurricane formation on tidally locked planets. Hurricanes can enhance the ocean mixing and
199 oceanic heat transport from warmer to cooler regions in both horizontal and vertical directions.
200 Hurricanes can also influence the transmission spectra of tidally locked planets. For instance, if
201 hurricane moves to the terminator, water vapor concentration would increase largely (Fig. S7),
202 which can influence the transmission signals. Unfortunately, present-day telescopes are not able to
203 observe this (39,40) mainly due to the small-scale height of the atmosphere and the relatively small
204 size of the hurricane to planetary radius. Differentiating them requires future large space telescopes
205 or ground-based extremely large telescopes.

206

Materials and Methods

This section includes experimental design, hurricane detection method, and the calculations of genesis potential index and Rossby deformation radius.

Experimental Design. The model we used is the global Community Atmosphere Model version 4 (CAM4) with a dynamical core of finite volume (41). Deep convection is parameterized using the updated mass flux scheme of Zhang and McFarlane (42). Subgrid-scale momentum transport associated with convection is included (43). The parameterization of shallow moist convection is based on Hack (44). Condensation, evaporation, and precipitation parameterization is based on Zhang et al. (45) and Rasch and Kristjansson (46). Cloud fraction is diagnosed from atmospheric stratification, convective mass flux, and relative humidity. Realistic radiative transfer of water vapor, clouds, greenhouse gases, and aerosols are included.

The horizontal resolution we employed is $0.47^{\circ} \times 0.63^{\circ}$ in latitude and longitude, respectively. The number of vertical levels is 26. Planetary surface is covered by seawater everywhere (an aquaplanet). Because of the high resolution and the limited computational power, we specify surface temperature in the simulations. The surface temperature is set according to previous simulations of lower-resolution AGCMs coupled to a 50-m slab ocean (47,48). On the day side, the surface temperature is a function of latitude and longitude: $(T_{max} - T_{min})\cos(\varphi)\cos(\lambda) + T_{min}$ (shown in Fig. 2A), or $(T_{max} - T_{min})\cos^{1/4}(\varphi)\cos^{1/4}(\lambda) + T_{min}$ (shown in Fig. S3), where T_{max} is the maximum surface temperature, T_{min} is the minimum surface temperature, φ is the latitude, and λ is the longitude. On the night side, the surface temperature is uniform with a value of T_{min} . Three groups of T_{max} and T_{min} are used. One is for planets near the inner edge of the habitable zone, 315 K & 310 K. The other two are for planets in the middle range of the habitable zone, 308 K & 275 K and 301 K & 268 K; the power of 1/4 is used due to the very weak temperature gradients in the substellar region and strong temperature gradients near the terminators.

232 Planetary rotation period is set to be equal to orbital period, i.e., 1:1 spin-orbit resonance. Four
233 rotation periods are examined: 6, 10, 20, and 40 earth days. For other types of spin-orbit resonance,
234 such as 3:2 like Mercury, the climate is in the condition between the synchronous rotation and the
235 very rapidly rotation of Earth (49), so that we have not done this kind of experiments. Planetary
236 radius and gravity are set to be the same as Earth, but both obliquity and eccentricity are set to zero.
237 Stellar temperature is set to 2600 or 3700 K; sensitivity test shows this has no significant effect on
238 the results. The stellar radiation at the substellar point is set to 1300 or 1800 W m⁻²; this also has
239 no significant effect on the results because the surface temperature is fixed. By default, the mean
240 surface pressure is 1 bar with ≈79% N₂ and ≈21% O₂. For greenhouse gases, we set CO₂
241 concentration to 367 parts per million by volume (ppmv), N₂O to 316 parts per billion by volume
242 (ppbv), and CH₄ to 1760 ppbv. Ozone concentration is set to be the same as present-day Earth,
243 which may influence the outflow temperature of hurricanes and the overshooting of extremely
244 strong convection.

245 In order to briefly test the effect of atmospheric composition on the hurricane formation, we
246 did several ideal experiments in which the background gas is set to H₂, He, N₂, O₂, and CO₂,
247 respectively. The corresponding mean molecular weights are 2.02, 4.00, 28.01, 31.99, and 44.00 g
248 mole⁻¹ and the corresponding specific heats (50) are 28.9, 20.8, 29.1, 29.5, and 37.2 J mole⁻¹ K⁻¹.
249 We modify these two constants only. The model we employed (as well as most other global models)
250 is unable to calculate the radiative transfer of dense H₂, He, O₂, or CO₂ and meanwhile surface
251 temperatures under background gases differing from Earth have not been seriously examined. So,
252 we choose to use globally uniform surface temperature (301 K) and stellar radiation (340 W m⁻²)
253 with neither seasonal nor diurnal cycle. Of course, this idealized thermal boundary condition is
254 unrealistic, but it is able to avoid the effects of strong wind shear, baroclinic zone, cold surface, or
255 other features that may inhibit hurricane formation or propagation (51). Two planetary rotation
256 periods are used, one and three earth days.

257 Initial states of the experiments were from long-term (40 earth years) simulations using a lower
258 resolution of $4^{\circ} \times 5^{\circ}$ or $1.9^{\circ} \times 2.5^{\circ}$ under the same experimental designs and parameterization
259 schemes. Then, each experiment was run for five earth years under the high resolution, and the last
260 four years were used to analyze below. As shown in Fig. 4, the southern hemisphere always has
261 more hurricanes than the northern hemisphere; this may be due to some asymmetry in initial state
262 or some stochastic process in the parameterizations.

263 Future work is required using AGCMs coupled to a slab ocean or using fully coupled
264 atmosphere–ocean models. Moreover, future work is required to examine how continents influence
265 the results. Hurricanes always decay quickly when they move over land because of the dramatic
266 reduction in evaporation and the increase in surface roughness. Global climate models with more
267 realistic cloud schemes and regional cloud resolving models with more accurate radiation transfer
268 are required to simulate the hurricane genesis, especially for those who have quite different
269 atmospheric compositions and/or air masses from Earth.

270
271 ***Hurricane detection.*** Hurricane formation and tracking is based on six hourly model output
272 variables using the Geophysical Fluid Dynamics Laboratory tracking algorithm (52). Candidate
273 hurricanes are identified by finding regions satisfying the following criteria: 1) the local 850-hPa
274 relative vorticity maximum exceeds $3.5 \times 10^{-5} \text{ s}^{-1}$; 2) the 850-hPa warm-core temperature must be
275 at least 0.5 K warmer than the surrounding local mean; 3) the distance between the local sea level
276 pressure minimum and the vorticity maximum should be within a distance of 2° latitude or
277 longitude, and so is the distance between the local sea level pressure minimum and the warm-core
278 center; 4) the maximum 850-hPa wind speed exceeds $\approx 33 \text{ m s}^{-1}$ at some time. These values of the
279 thresholds influence the exact number of detected hurricanes but do not affect the main conclusions
280 of this study.

282 **Genesis Potential Index (GPI).** The GPI is written as:

$$283 \quad GPI = |10^5(\zeta + f)|^{3/2}(RH/50)^3(V_{pot}/70)^3(1.0 + 0.1 * V_{shear})^{-2}, \quad (1)$$

284 where ζ is the vertical component of relative vorticity, f is the planetary vorticity, RH is the
285 relative humidity at the middle troposphere (600 hPa), V_{shear} is the wind shear of horizontal winds
286 between the upper and lower troposphere (300 minus 850 hPa; or called vertical wind shear), and
287 V_{pot} is potential intensity. V_{pot} is a measure of the maximum near-surface wind that can be
288 maintained by hurricane under given environmental condition. Note that these parameters are not
289 all independent; for instance, vertical wind shear can influence relative humidity.

290 The value of V_{pot} is calculated based on a local balance between thermal energy import and
291 mechanical energy dissipation (53,54), written as

$$292 \quad V_{pot} = \frac{C_k T_s - T_o}{C_D T_o} [CAPE^* - CAPE^b] |_{m}, \quad (2)$$

293 where C_k is the exchange coefficient for enthalpy, C_D is the drag coefficient, T_s is the sea surface
294 temperature, T_o is the mean outflow temperature, $CAPE^*$ is the convective available potential
295 energy of air lifted from saturation at sea surface, and $CAPE^b$ is that of the boundary layer air.
296 Both $CAPE^*$ and $CAPE^b$ are computed at the radius of maximum surface wind.

297
298 **Rossby deformation radius.** The Rossby radius of deformation (L_R) is the length scale at which
299 rotational effects become as important as the effects of gravity waves or buoyancy in the evolution
300 of the flow in a disturbance. L_R is equal to $\frac{NH}{f}$, where N is the Brunt-Vaisala frequency, and f is

301 the Coriolis parameter. H is the scale height, equaling to $\frac{R^*\bar{T}}{M_d g}$, where R^* is the universal gas
302 constant, \bar{T} is the mean air temperature, M_d is the molar weight of the atmosphere, and g is the
303 gravity (55). In idealized experiments with uniform surface temperature or uniform rotation, it is
304 one of the rough scales that can be used for understanding hurricane size (56-58). In our experiments

305 shown in Fig. 5, the value of L_R is ≈ 500 -1500 km, comparable to the hurricane size. However, L_R

306 is strongly latitude-dependent because f is equal to $2\Omega\sin(\varphi)$ (Ω is the rotation rate and φ is the
307 latitude), but the hurricane size does not exhibit the same dependence. A better scaling was not
308 found because of the nonlinear dynamics of hurricane and the complex interactions between
309 hurricane and diabatic heating, environmental relative humidity, mesoscale convective system, and
310 other features.

311

312 **References**

- 313 1. R. A. Anthes, *Tropical cyclones—their evolution, structure and effects*. (American Meteor
314 Society, Boston, MA, 1982).
- 315 2. K. A. Emanuel, 100 Years of Progress in Tropical Cyclone Research. *Meteor. Monogr.* **59**,
316 15.1-15.68 (2019).
- 317 3. K. A. Emanuel, Contribution of tropical cyclones to meridional heat transport by the oceans.
318 *J. Geophys. Res.* **106**, 14771-14781 (2001).
- 319 4. M. Jansen, R. Ferrari, Impact of the latitudinal distribution of tropical cyclones on ocean heat
320 transport. *Geophys. Res. Lett.* **36**, 150-164 (2009).
- 321 5. H. Li, R. L. Sriver, Impact of Tropical Cyclones on the Global Ocean: Results from
322 Multidecadal Global Ocean Simulations Isolating Tropical Cyclone Forcing. *J. Clim.* **31**, 8761-
323 8784 (2018).
- 324 6. A. P. Showman, R. D. Wordsworth, T. M. Merlis, Y. Kaspi, *Atmospheric Circulation of*
325 *Terrestrial Exoplanets*. (Univ Ariz Press, Tucson, AZ, 2013).
- 326 7. R. T. Pierrehumbert, M. Hammond, Atmospheric Circulation of Tide-Locked Exoplanets.
327 *Annu. Rev. Fluid Mech.* **51**, 275-303 (2019).
- 328 8. R. K. Kopparapu et al. Habitable Moist Atmospheres on Terrestrial Planets near the Inner Edge
329 of the Habitable Zone around M Dwarfs. *Astrophys. J.* **845**, 5 (2017).

- 330 9. J. Haqq-Misra, E. T. Wolf, M. Joshi, X. Zhang, R. K. Kopparapu, Demarcating circulation
331 regimes of synchronously rotating terrestrial planets near the inner edge of the habitable zone.
332 *Astrophys. J.* **852** (2017).
- 333 10. J. Yang, D. S. Abbot, D. B. Koll, Y. Hu, A. P. Showman, Ocean dynamics and the inner edge
334 of the habitable zone for tidally locked terrestrial planets. *Astrophys. J.* **871**, 1-17 (2019).
- 335 11. A. D. Del Genio et al. Habitable climate scenarios for Proxima Centauri b with a dynamic
336 ocean. *Astrobiology* **19**, 99-125 (2019).
- 337 12. M. Turbet et al. Modeling climate diversity, tidal dynamics and the fate of volatiles on
338 TRAPPIST-1 planets. *Astron. Astrophys.* **612**, 1-22 (2018).
- 339 13. L. Carone, R. Keppens, L. Decin, Connecting the dots--II. Phase changes in the climate
340 dynamics of tidally locked terrestrial exoplanets. *Mon. Not. R. Astron. Soc.* **453**, 2412-2437
341 (2015).
- 342 14. R. Wordsworth, Atmospheric heat redistribution and collapse on tidally locked rocky planets.
343 *Astrophys. J.* **806**, 180 (2015).
- 344 15. A. L. Shields, S. Ballard, J. A. Johnson, The habitability of planets orbiting M-dwarf stars.
345 *Phys. Rep.* **663**, 1-38 (2016).
- 346 16. M. M. Joshi, R. M. Haberle, R. T. Reynolds, Simulations of the Atmospheres of Synchronously
347 Rotating Terrestrial Planets Orbiting M Dwarfs: Conditions for Atmospheric Collapse and the
348 Implications for Habitability. *Icarus* **129**, 450-465 (1997).
- 349 17. T. M. Merlis, T. Schneider, Atmospheric dynamics of Earth-like tidally locked aqua-planets.
350 *J. Adv. Model Earth Syst.* **2**, 13 (2010).
- 351 18. A. Edson, S. Lee, P. Bannon, J. F. Kasting, D. Pollard, Atmospheric circulations of terrestrial
352 planets orbiting low-mass stars. *Icarus* **212**, 1-13 (2011).
- 353 19. R. T. Pierrehumbert, A palette of climates for Gliese 581g. *Astrophys. J. Lett.* **726**, L8-L12
354 (2011).

- 355 20. J. Leconte et al. 3D climate modeling of close-in land planets: circulation patterns, climate
356 moist bistability and habitability. *Astron. Astrophys.* **554**, A19(2013).
- 357 21. K. Menou, Water-trapped worlds. *Astrophys. J.* **774**, 51 (2013).
- 358 22. S. Noda et al. The circulation pattern and day-night heat transport in the atmosphere of a
359 synchronously rotating aquaplanet: dependence on planetary rotation rate. *Icarus* **282**, 1-18
360 (2017).
- 361 23. J. Checlair, K. Menou, D. S. Abbot, No snowball on habitable tidally locked planets. *Astrophys.*
362 *J.* **845**, 132 (2017).
- 363 24. I. A. Boutle et al. Exploring the climate of Proxima B with the Met Office Unified Model.
364 *Astron. Astrophys.* **601**, A120 (2017).
- 365 25. J. Bin, F. Tian, Y. Lin, Y. Wang, Low Probability of Tropical Cyclones on Ocean Planets in
366 the Habitable Zones of M Dwarfs. *Icarus* **299**, 364-369 (2017).
- 367 26. J. G. Charney, A. Eliassen, On the growth of the hurricane depression. *J. Atmos. Sci.* **21**, 68-
368 75 (1964).
- 369 27. R. K. Smith, On the theory of CISK. *Q. J. R. Meteorol. Soc.* **123**, 407-418 (1997).
- 370 28. Y. Q. Wang, Hurricane dynamics, in Tropical cyclones and hurricanes, *In: Zhang GRNP (ed)*
371 *Encyclopedia of atmospheric sciences, 2nd edn.* (Academic Press, Oxford, 2015), pp. 8-29.
- 372 29. M. Yamasaki, A View on Tropical Cyclones as CISK. *J. Meteorol. Soc. Jpn.* **85B**, 145-164
373 (2007).
- 374 30. A. E. Gill, Some simple solutions for heat-induced tropical circulation. *Q. J. R. Meteorol. Soc.*
375 **106**, 447-462 (1980).
- 376 31. A. P. Showman, L. M. Polvani, Equatorial superrotation on tidally locked exoplanets.
377 *Astrophys. J.* **738**, 71 (2011).
- 378 32. T. Matsuno, Quasi-Geostrophic Motions in the Equatorial Area. *J. Meteorol. Soc. Jpn.* **44**, 25-
379 42 (1966).

- 380 33. K. A. Emanuel, D. S. Nolan, Tropical cyclone activity and global climate. Preprints, *26th*
381 *Conference on Hurricanes and Tropical Meteorology*, Miami, FL, Amer. Meteor. Soc. 240-
382 241 (2004).
- 383 34. S. J. Camargo, K. A. Emanuel, A. H. Sobel, Use of a Genesis Potential Index to Diagnose
384 ENSO Effects on Tropical Cyclone Genesis. *J. Clim.* **20**, 4819-4834 (2007).
- 385 35. B. Tang, K. A. Emanuel, Midlevel Ventilation's Constraint on Tropical Cyclone Intensity. *J.*
386 *Atmos. Sci.* **67**, 1817-1830 (2010).
- 387 36. C. Li, A. P. Ingersoll, Moist convection in hydrogen atmospheres and the frequency of Saturn's
388 giant storms. *Nat. Geosci.* **8**, 398-403 (2015).
- 389 37. J. Leconte, F. Selsis, F. Hersant, T. Guillot, Condensation-inhibited convection in hydrogen-
390 rich atmospheres. *Astron. Astrophys.* **598**, A98 (2016).
- 391 38. T. Guillot, Condensation of methane, ammonia, and water and the inhibition of convection in
392 giant planets. *Science* **269**, 1697-1699 (1995).
- 393 39. C. V. Morley, L. Kreidberg, Z. Rustamkulov, T. Robinson, J. J. Fortney, Observing the
394 atmospheres of known temperate earth-sized planets with JWST. *Astrophys. J.* **850**, 121
395 (2017).
- 396 40. J. de Wit et al. Atmospheric reconnaissance of the habitable-zone Earth-sized planets orbiting
397 TRAPPIST-1. *Nat. Astron.* **2**, 214-219 (2018).
- 398 41. R. Neale et al. *Description of the NCAR Community Atmosphere Model (CAM 4.0)*. NCAR
399 *Technical Note NCAR/TN-464+STR* (National Center for Atmospheric Research, Boulder, CO,
400 2010).
- 401 42. G. J. Zhang, N. A. McFarlane, Sensitivity of climate simulations to the parameterization of
402 cumulus convection in the Canadian Climate Centre general circulation model. *Atmos.-Ocean*
403 **33**, 407-446 (1995).

- 404 43. J. H. Richter, P. J. Rasch, Effects of Convective Momentum Transport on the Atmospheric
405 Circulation in the Community Atmosphere Model, Version 3. *J. Clim.* **21**, 1487-1499 (2008).
- 406 44. J. J. Hack, Parameterization of moist convection in the National Center for Atmospheric
407 Research community climate model (CCM2). *J. Geophys. Res.* **99**, 5551-5568 (1994).
- 408 45. M. Zhang, W. Lin, C. S. Bretherton, J. J. Hack, P. J. Rasch, A modified formulation of
409 fractional stratiform condensation rate in the NCAR Community Atmospheric Model (CAM2).
410 *J. Geophys. Res.* **108**, 4035 (2003).
- 411 46. P. J. Rasch, J. E. Kristjánsson, A Comparison of the CCM3 Model Climate Using Diagnosed
412 and Predicted Condensate Parameterizations. *J. Clim.* **11**, 1587-1614 (1998).
- 413 47. J. Yang, N. B. Cowan, D. S. Abbot, Stabilizing cloud feedback dramatically expands the
414 habitable zone of tidally locked planets. *Astrophys. J. Lett.* **771**, L45 (2013).
- 415 48. E. T. Wolf, R. K. Kopparapu, J. Haqq-Misra, Simulated Phase-dependent Spectra of Terrestrial
416 Aquaplanets in M Dwarf Systems. *Astrophys. J.* **877**, 35 (2019).
- 417 49. J. Yang, G. Boue, D. C. Frabrycky, D. S. Abbot, Strong Dependence of the Inner Edge of the
418 Habitable Zone on Planetary Rotation Rate. *Astrophys. J. Lett.* **787**, L2 (2014).
- 419 50. X. Zhang, A. P. Showman, Effects of Bulk Composition on the Atmospheric Dynamics on
420 Close-in Exoplanets. *Astrophys. J.* **836**, 73 (2017).
- 421 51. T. M. Merlis, I. M. Held, Aquaplanet Simulations of Tropical Cyclones. *Curr. Clim. Change.*
422 *Rep.* **5**, 185-195 (2019).
- 423 52. M. Zhao, I. M. Held, S. J. Lin, G. A. Vecchi, Simulations of Global Hurricane Climatology,
424 Interannual Variability, and Response to Global Warming Using a 50-km Resolution GCM. *J.*
425 *Clim.* **22**, 6653-6678 (2009).
- 426 53. K. A. Emanuel, The maximum intensity of hurricanes. *J. Atmos. Sci.* **45**, 1143-1155 (1988).
- 427 54. M. Bister, K. A. Emanuel, Low frequency variability of tropical cyclone potential intensity 1.
428 Interannual to interdecadal variability. *J. Geophys. Res.* **107**, ACL 26-21-ACL 26-15 (2002).

429 55. J. M. Wallace, P. V. Hobbs, *Atmospheric Science: An Introductory Survey*. 2nd ed. (Academic
430 Press, 2006).

431 56. I. M. Held, M. Zhao, Horizontally homogeneous rotating radiative-convective equilibria at
432 GCM resolution. *J. Atmos. Sci.* **65**, 2003-2013 (2008).

433 57. D. R. Chavas, N. Lin, W. Dong, Y. Lin, Observed tropical cyclone size revisited. *J. Clim.* **29**,
434 2923-2939 (2016).

435 58. D. R. Chavas, K. A. Reed, Dynamical Aquaplanet Experiments with Uniform Thermal
436 Forcing: System Dynamics and Implications for Tropical Cyclone Genesis and Size. *J. Atmos.*
437 *Sci.* **76**, 2257-2274 (2019).

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447 **materials availability:** The data that support the plots within this paper will be stored on a public
448 structured data repository after the paper accepted. The source codes of the model can be
449 downloaded from <http://www.cesm.ucar.edu/models/>, and changes of the model are available from
450 the authors J.Y. junyang@pku.edu.cn and M.Y. yanmy@pku.edu.cn.

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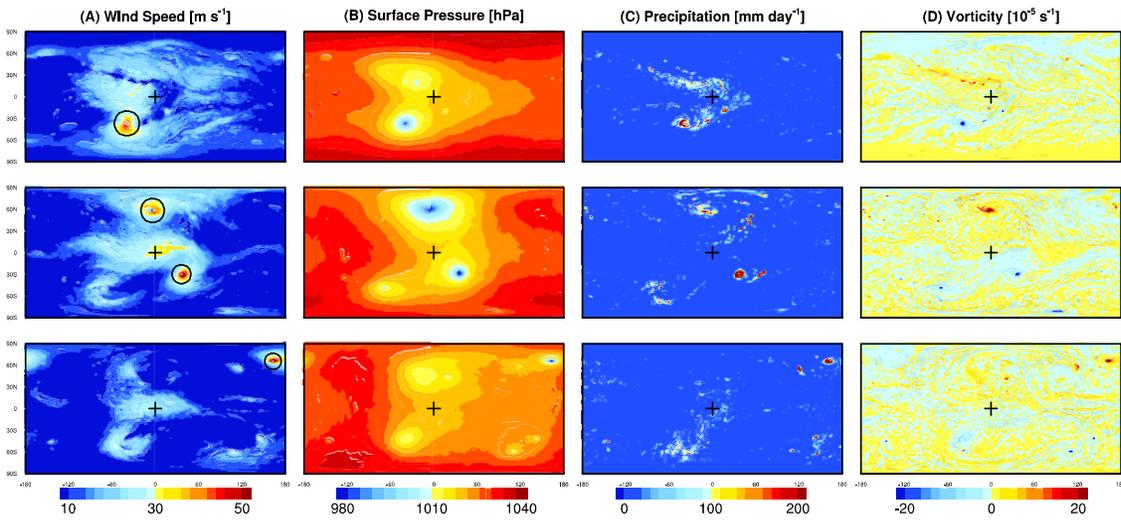


Fig. 1. Snapshots of hurricanes on a tidally locked aqua-planet near the inner edge of the habitable zone in the control experiment. From left to right, the variables are instantaneous wind speed at 850 hPa, surface air pressure, precipitation, and the vertical component of relative vorticity at 850 hPa, respectively. From upper to bottom, they are for three different moments. The four hurricanes are marked with black circles over the wind speed panels. The black cross is the substellar point in this figure and hereafter. See Video S1 online for the evolution of the hurricanes.

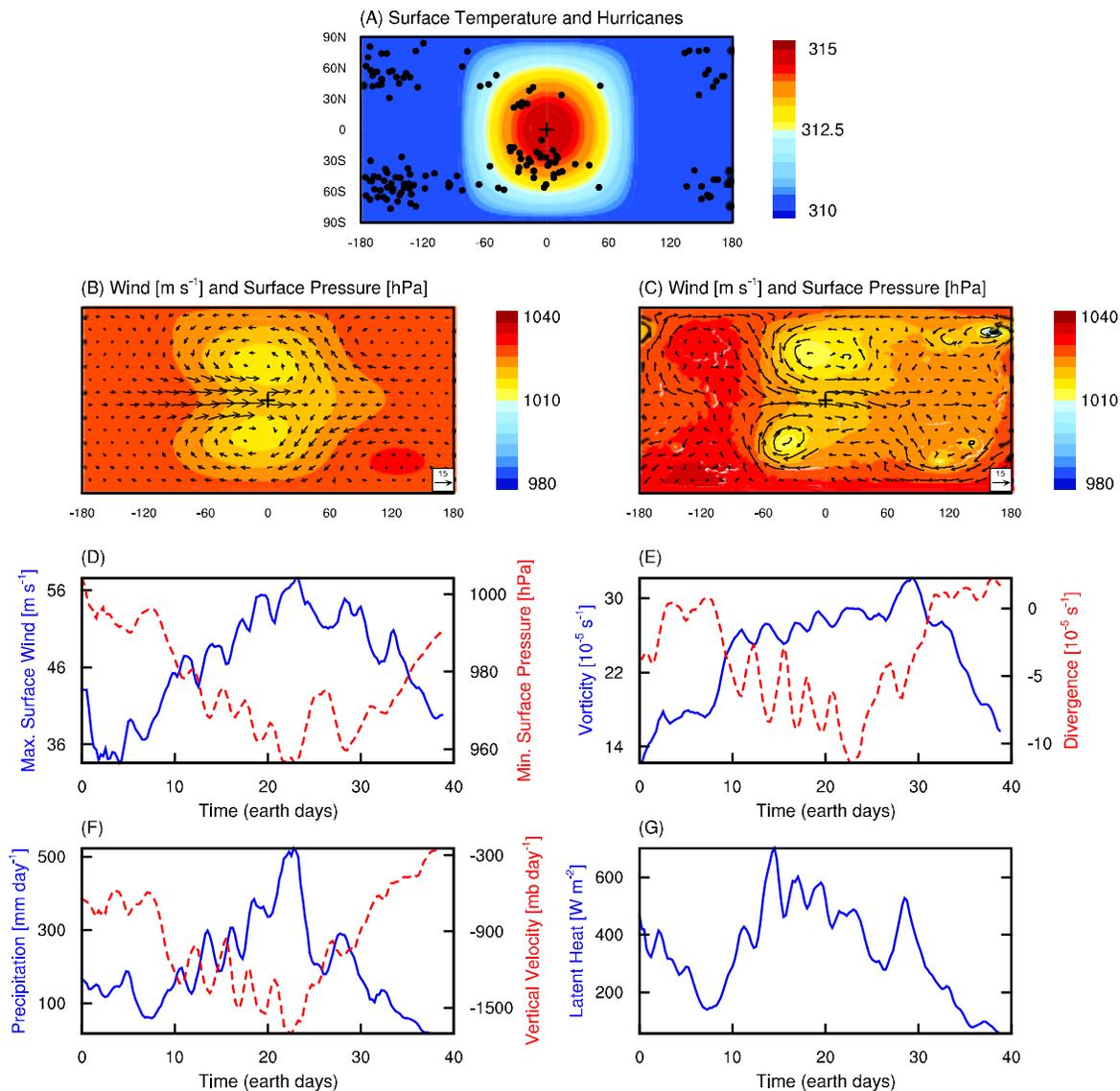
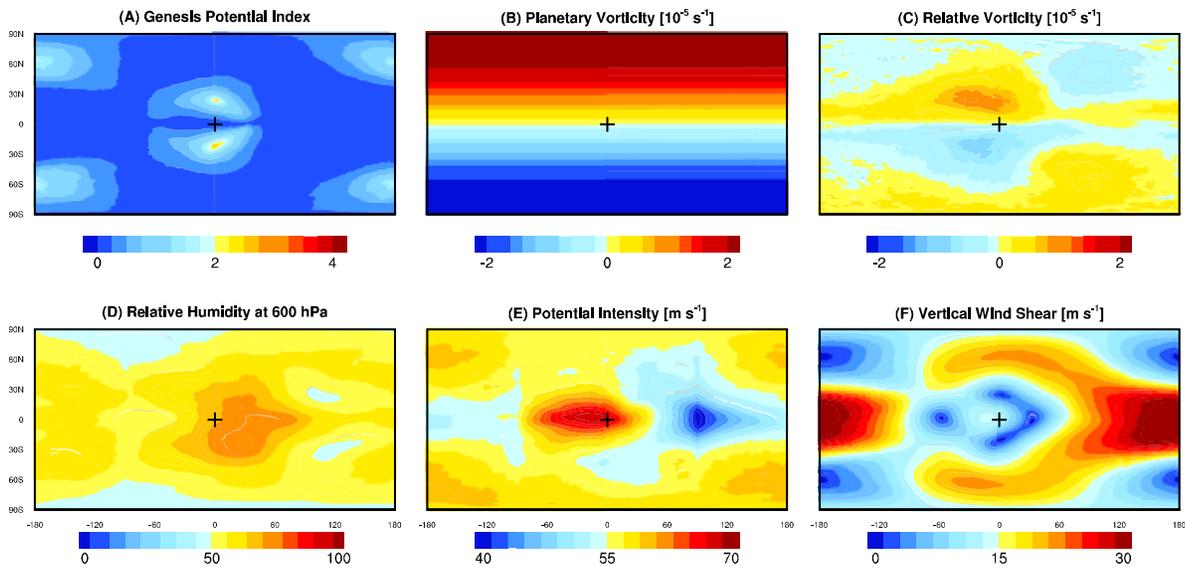


Fig. 2. Mechanisms for hurricane formation in the control experiment. (A): The location of hurricane formation (dots) and surface air temperature (color shading). The number of hurricanes during the four-years integration is 154. (B): Long-term mean surface air pressure (shading) and winds at 850 hPa (vector). (C): Same as (B) but for an instantaneous. The life cycle of one hurricane on the day side: (D): maximum surface wind speed (blue) and minimum surface pressure (red), (E): relative vorticity (blue) and divergence (red) at 850 hPa, (F): precipitation (blue) and vertical velocity at 850 hPa (red), and (G): surface latent heat flux (blue). For (E), (F) and (G), the variables are calculated for area mean of $500 \times 500 \text{ km}^2$ around the low-pressure center.



472
 473 **Fig. 3. Environmental condition for hurricane formation.** (A): Genesis potential index (GPI);
 474 (B): planetary vorticity; (C): long-term mean relative vorticity at 850 hPa; (D): relative humidity at
 475 600 hPa; (E): potential intensity; (F): vertical shear of the horizontal winds between 300 and 850
 476 hPa. Note the environmental vorticity in panel (C) is one order smaller than the vorticity of
 477 hurricanes shown in Fig. 1D.

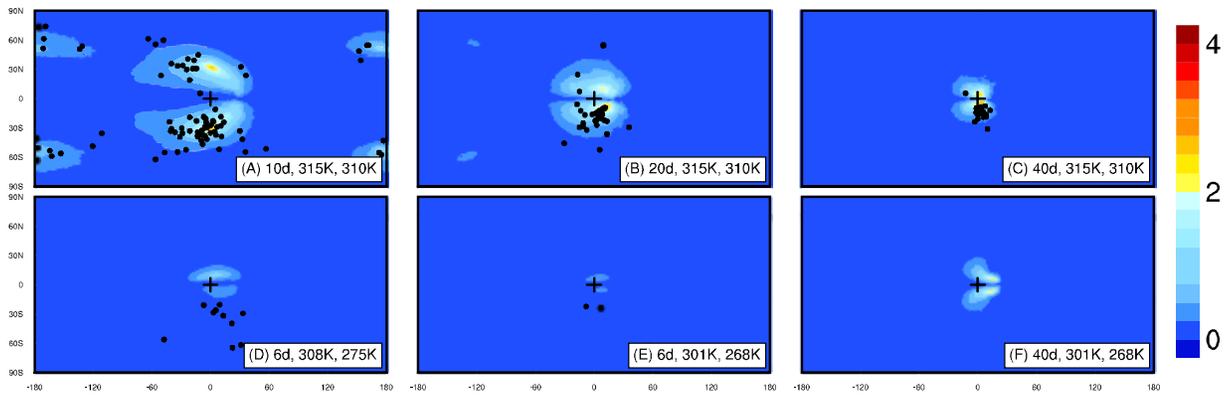
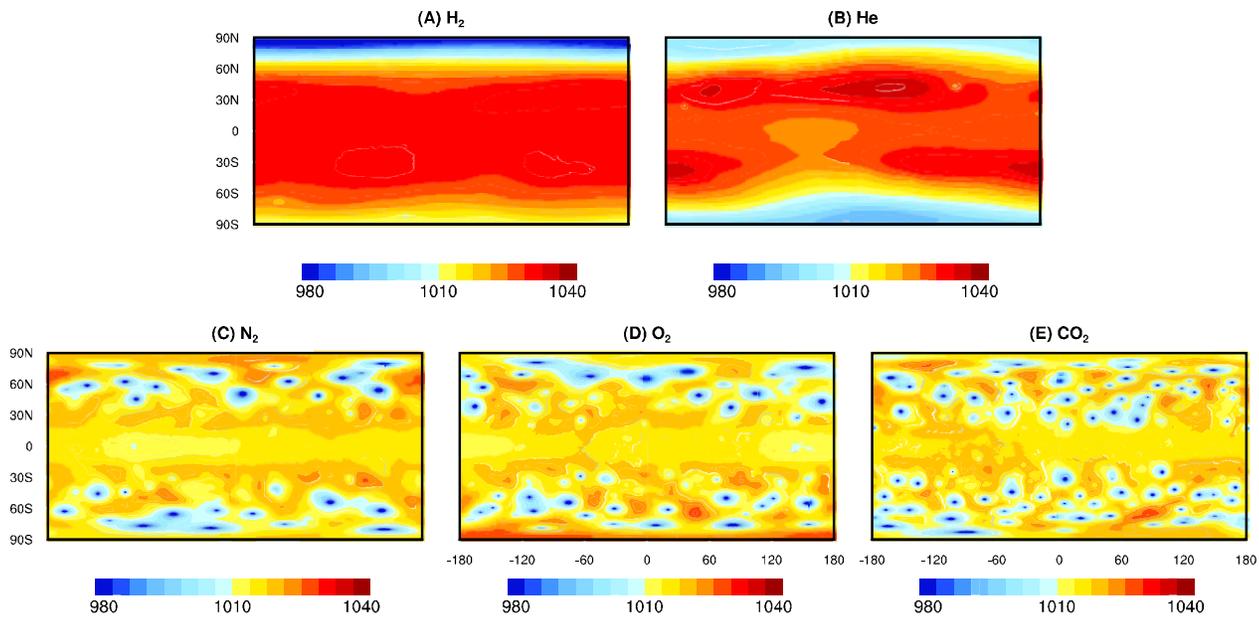


Fig. 4. Effects of planetary rotation and surface temperature on hurricane formation (dots) and the GPI (color shading). Experiments are for varying rotation period (**A-C**), varying surface temperature (**D-E**), and varying both (**F**). Experimental designs are the same as the control experiment in Fig. 1 except that the rotation period is set to (**A**): 10, (**B**): 20, and (**C**): 40 earth days; (**D**): the maximum surface temperature is reduced from 315 to 308 K and the night-side surface temperature is reduced from 310 to 275 K (see Fig. S3A); (**E**): same as (**D**) but for 301 K and 268 K, respectively (see Fig. S3B); and (**F**) same as (**E**) but for a rotation period of 40 earth days. The number of hurricanes is 88, 34, 21, 10, 2, and 0, respectively. See Fig. S4 for the snapshots of typical hurricanes.



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Fig. 5. Effects of background gases on hurricane formation. From (A) to (E), they are snapshots of instantaneous surface air pressure (hPa) under background gases of H₂, He, N₂, O₂, and CO₂, respectively. In all these experiments, the planetary rotation period is 1 earth day and surface temperature is uniform (301 K). For experiments with a rotation period of 3 earth days, the results are the same except that the hurricanes are larger in size but less in number in the latter three experiments.